

The role of diffusion in nonequilibrium phase transitions

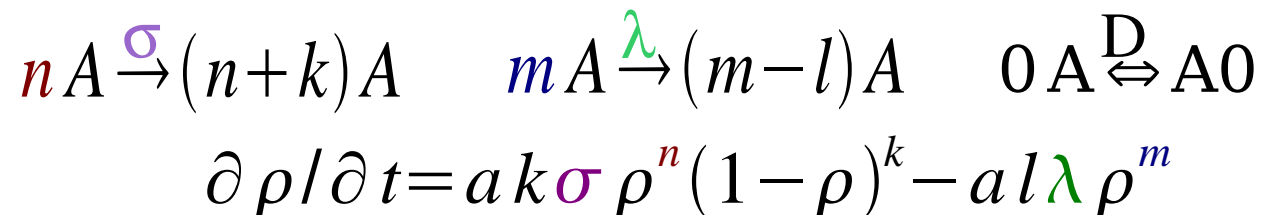
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- One of the primary goals is to explore universality classes like in equilibrium statistical physics
- Due to the lack of Gibbs distribution phase transitions (PT) may occur in low dimensions
- Which factors determine the PT universality class of a model of short range interactions ? Besides the spatial dimensions, boundaries, inhomogenities:
 - 1) Symmetries, conservation laws like in equilibrium (BAW2 ...)
 - 2) Initial conditions (temporal boundary condition) (PCP ...)
 - 3) Topological effects in low dimensions (multi-comp systems...)
 - 4) Dynamically generated long range memory (coupled systems...)
 - 5) Mean-field classes of RD : $nA \rightarrow (n+k)A$, $mA \rightarrow (m-1)A \rightarrow /$.
 - 6) For competing dynamics (diffusion) $\rightarrow /$.

See also: *G. Ódor, Rev. Mod. Phys. 76 (2004) no. 3.*

Mean-field classes of site restricted, one-component reaction-diffusion systems

- General, reaction-diffusion systems :



Order parameter ρ

$$\rho(t) \propto t^{-\alpha}$$

$$\rho(\infty) \propto \epsilon^\beta$$

- $n = m$: $\beta = 1, \alpha = 1/n$ $\sigma_c = l/(k+l)$
- $n < m$: $\beta = 1/(m-n), \alpha = 1/(m-1)$ $\sigma_c = 0$
- $n > m$: First order transition

n és m determine the mean-field class!

Diffusion does not play a role.

G. Ódor: PRE 67, 056114 (2003)

The dynamical cluster mean-field method

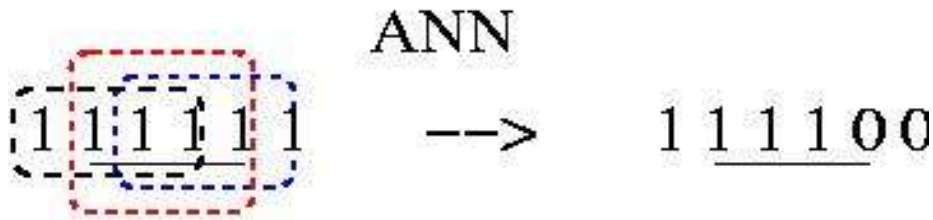
Master equations for $N=1,2,\dots$ block probabilities:

$$\partial P_N(s_i)/\partial t = f(P_N(s_i)), \quad s_i = 0, 1 \quad P_{abcde} \simeq P_{abcd} P_{bcde} / P_{bcd}$$

Example 4-Blocks:

Conf. 1

Conf. 2

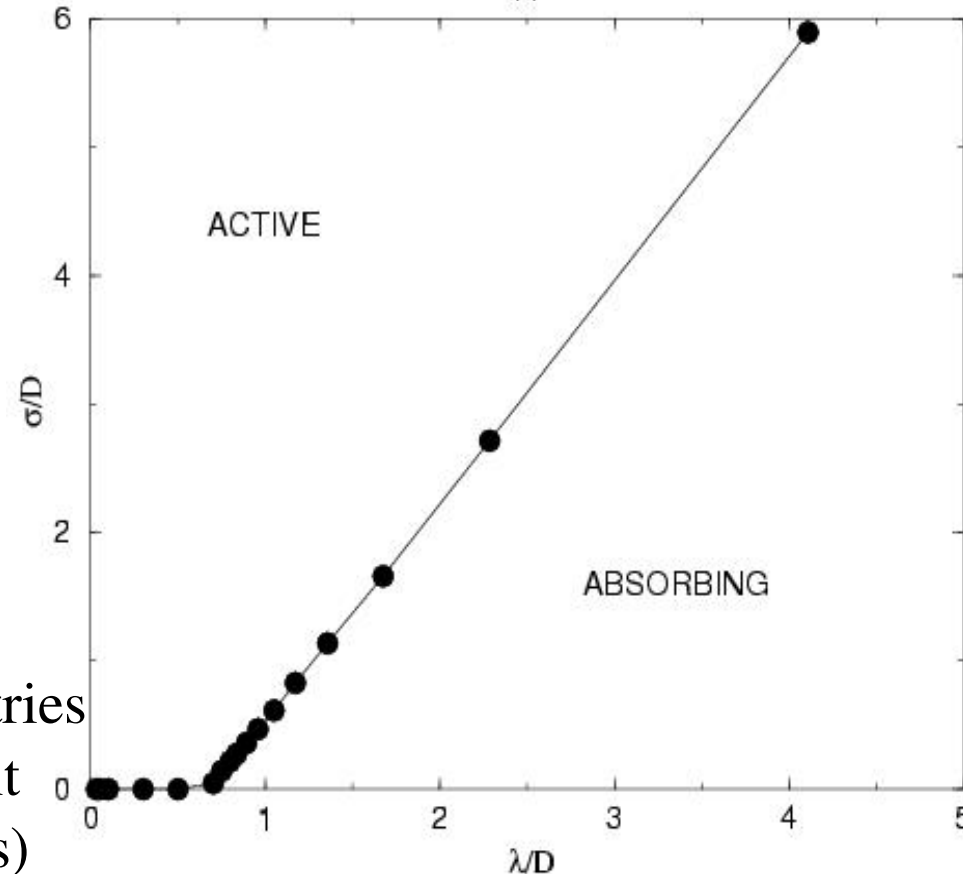


$$\lambda \frac{P_{1111} P_{1111} P_{1111}}{(P_{1110} + P_{1111})^2} P_{111}$$

Taking into account symmetries for $N=8$ case 136 independent variables (block. probabilities)

A \rightarrow 2A, 2A \rightarrow 0 model

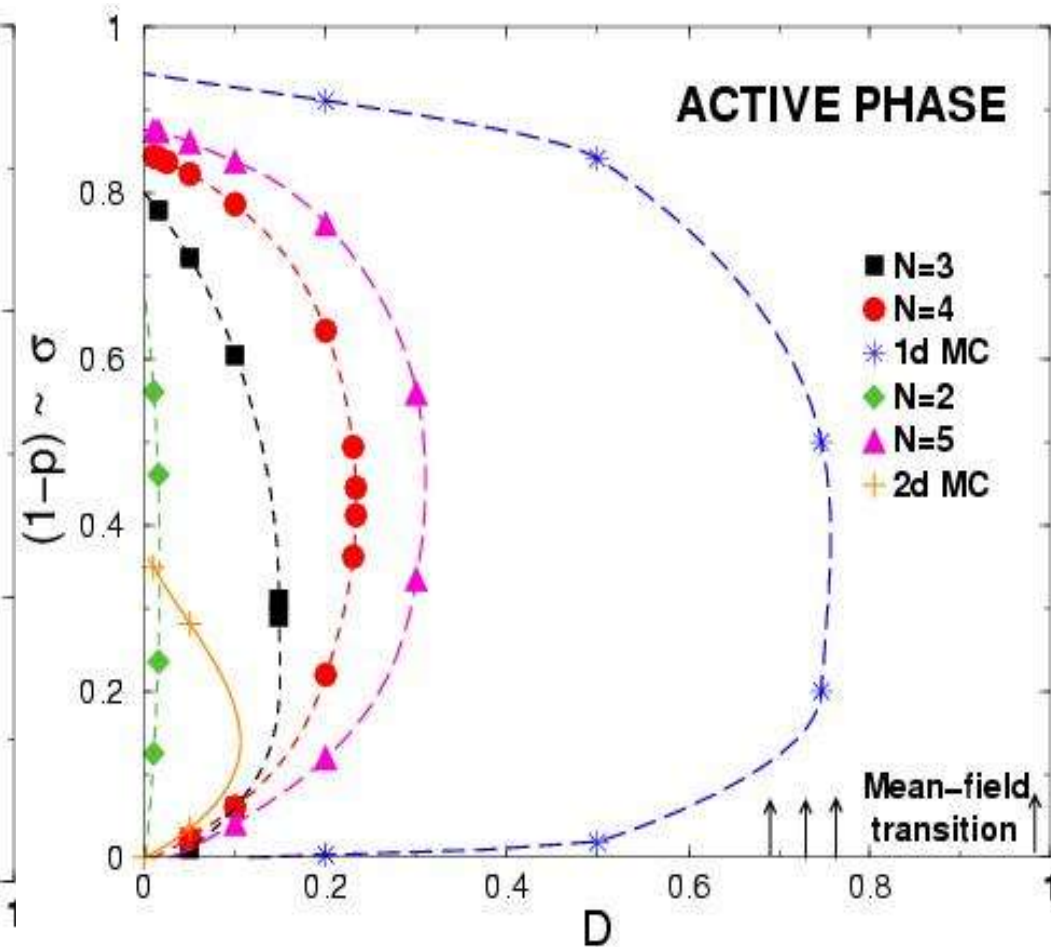
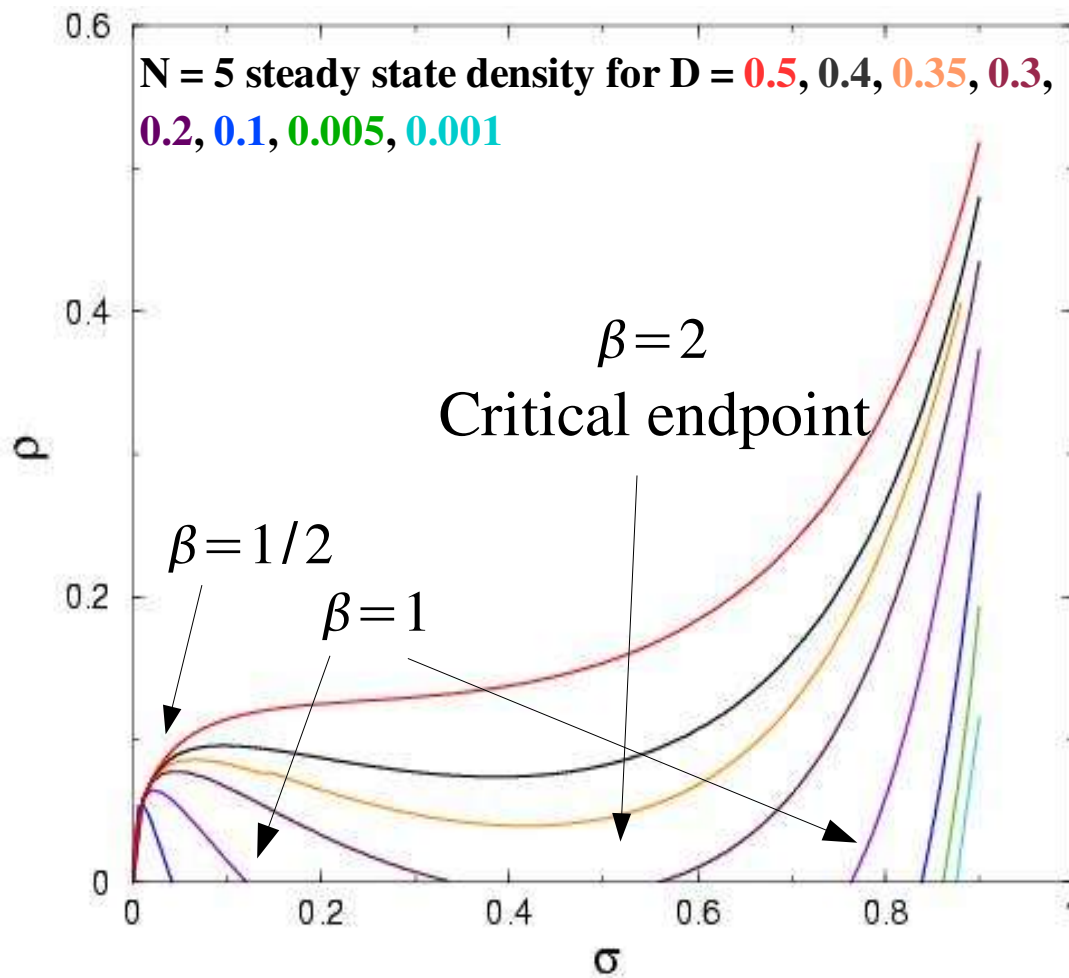
4-cluster approximation



Cluster approximations for: $2A \rightarrow 3A$, $4A \rightarrow 0$ model

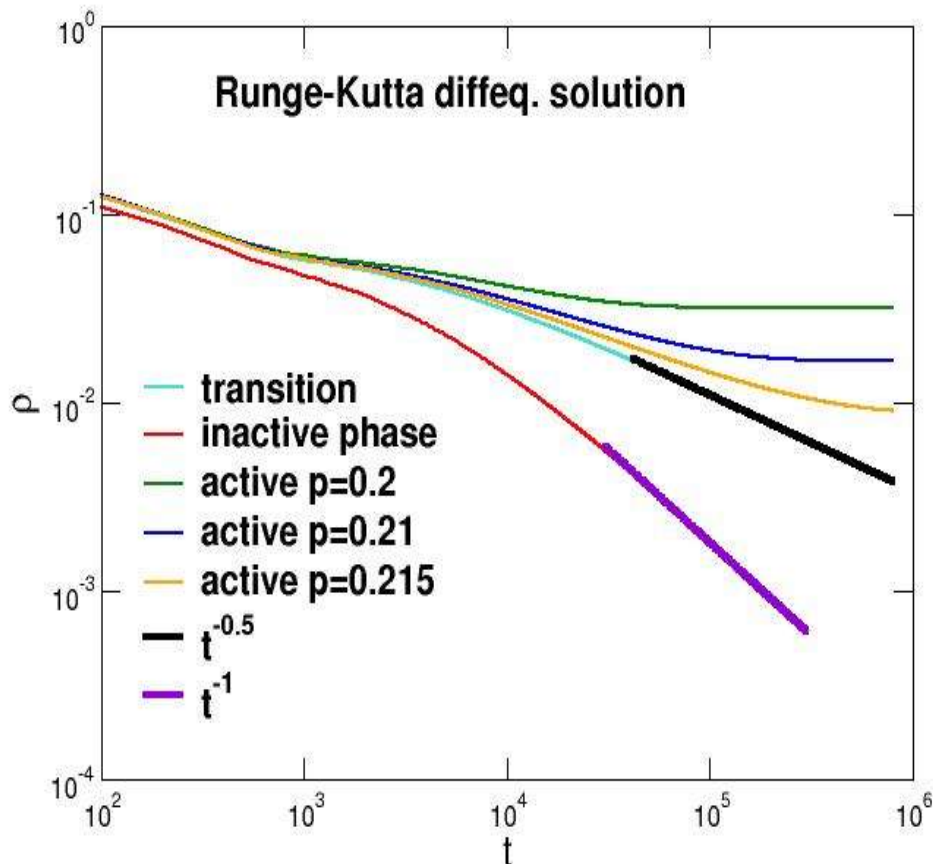
Steady state density for $N \geq 2$ (diffusion dependence).

Unexpected phase transitions for $\sigma_c > 0$, with $\beta=1$ (for $n < m$ MF: $\sigma_c = 0$)



Cluster approximations for: $2A \rightarrow 3A$, $4A \rightarrow 0$ model

Density decay solution for $N = 3$ at $D=0.05$ near $\sigma_c > 0$ critical point with exponent $\alpha = 0.5$ ($2A \rightarrow 3A$, $2A \rightarrow 0$ (PCPD) behavior)



Simulations in 1 and 2 dimension support N-cluster results:

G. Ódor, PRE 69, 036112 (2004)

G. Ódor: cond-mat/0403562

Unexpected by perturbative RG

For low diffusions the $2A \rightarrow 3A \rightarrow 4A \rightarrow 0$ process becomes relevant !

Lattice simulations

One dimension

production

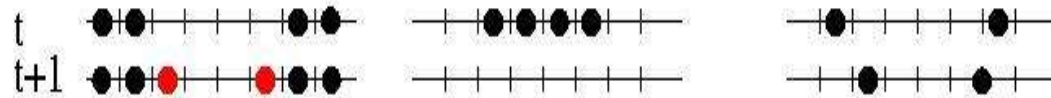
$$\sigma : (1-p)(1-D)/2$$

annihilation

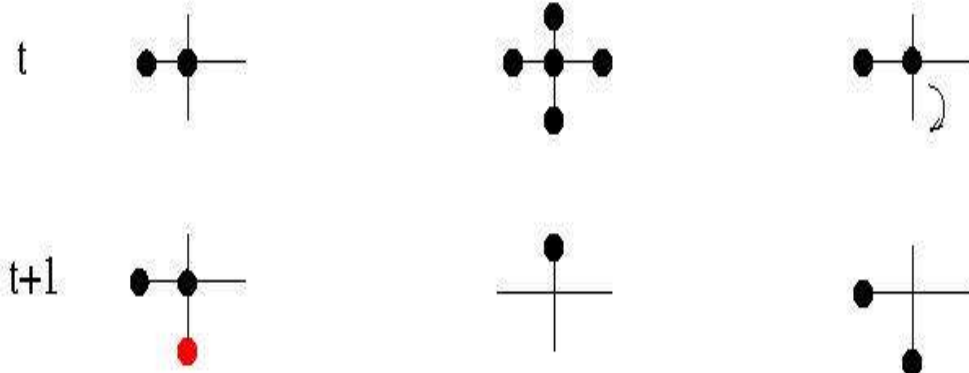
$$\lambda = p(1-D)$$

diffusion

rate: D



Two dimensions

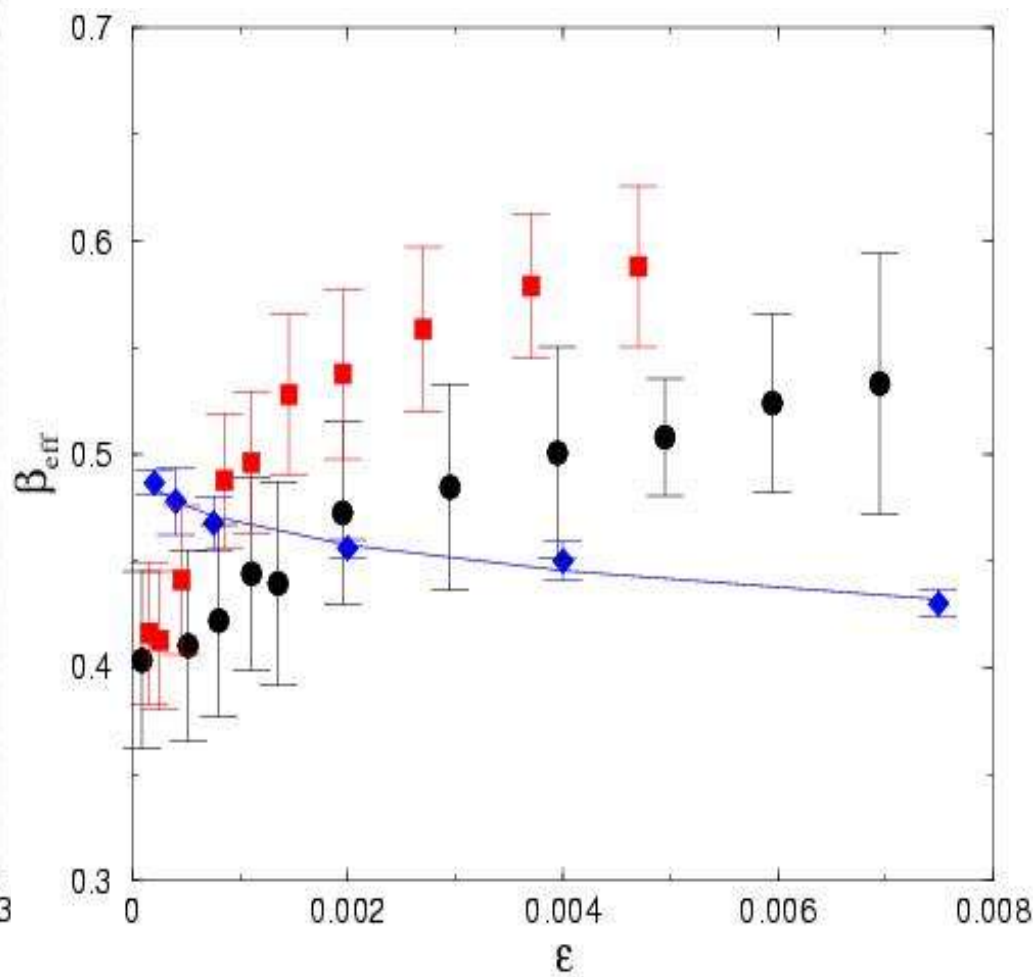
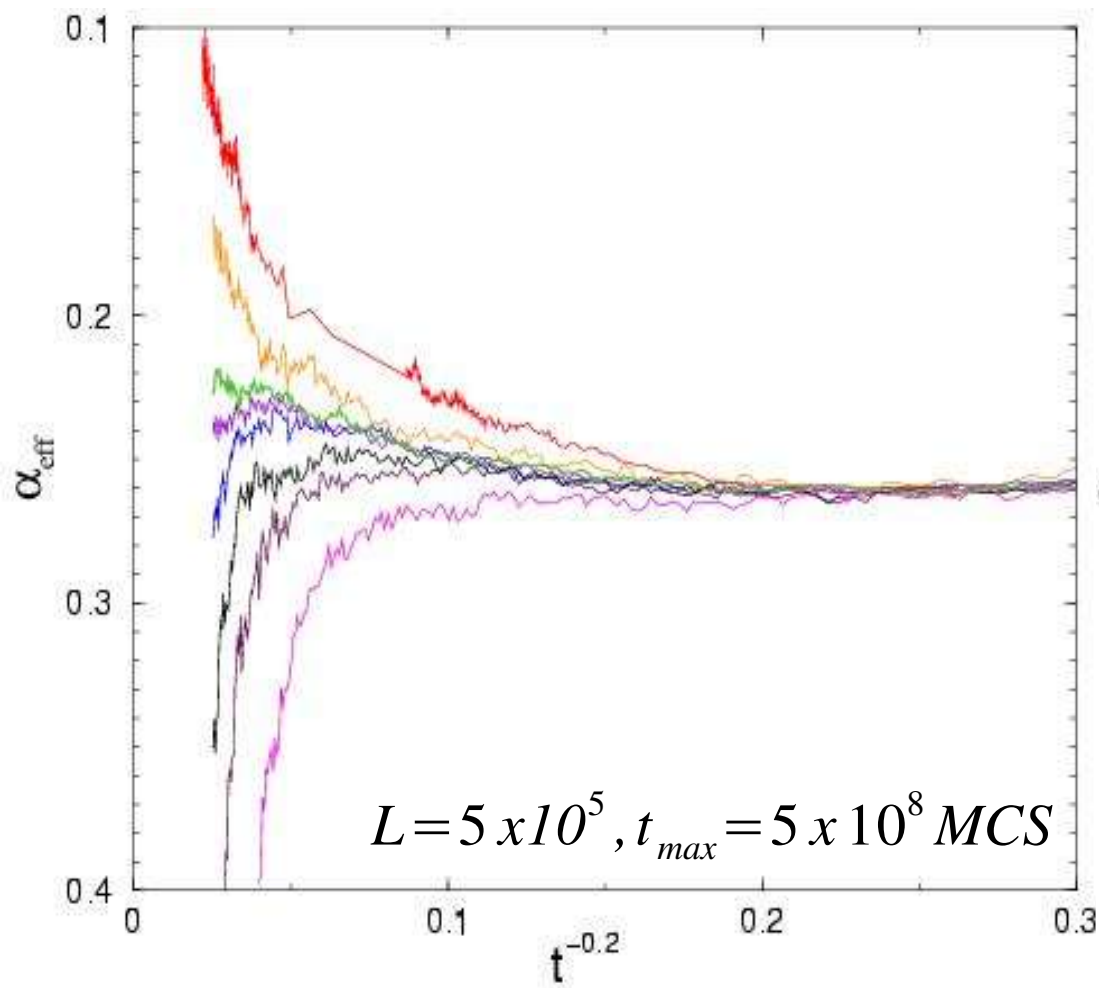


- Random sequential update of occupied sites
- Full or random half filled initial conditions
- $L \leq 500\,000$ in 1d,
 $L \leq 7000 \times 7000$ in 2d
periodic boundary conditions
- $t_{\max} \leq 5 \times 10^8$ MCS
- Master-slave parallel algorithms on international computing GRIDS ($\sim 100 - 500$ CPU-s)

Simulation results for the $2A \rightarrow 3A, 4A \rightarrow 0$ model in 1 dimension

Density decay local exponents
with : $\alpha = 0.21(2)$ (\sim PCPD), at
 $D = 0.5, \sigma_c = 0.42075$.

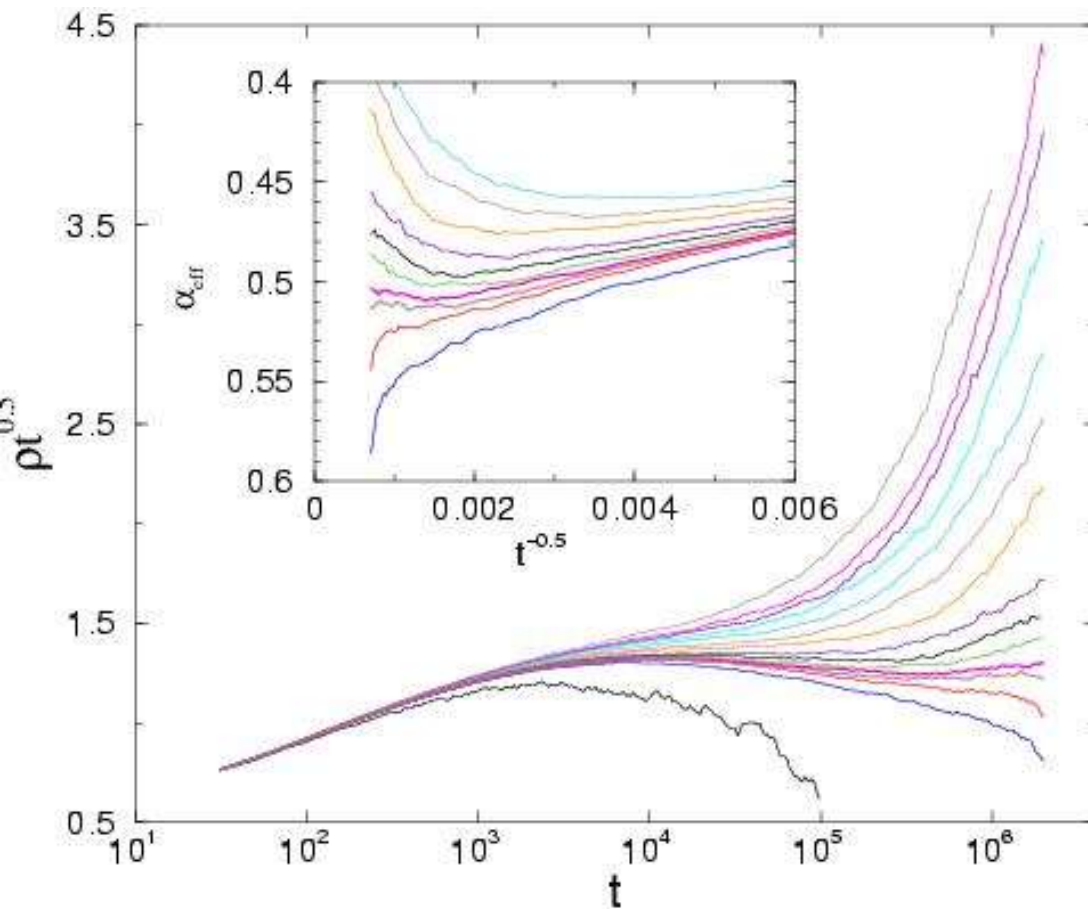
Steady state local exponents near $\sigma_c = 0.42$
for $D = 0.5, 0.2$, with $\beta = 0.40(2)$ (\sim PCPD).
For $D = 0.9$: only mean-field with $\beta = 1/2$.
 $2A \rightarrow 0$ process becomes irrelevant !



Simulation results for the $2A \rightarrow 3A, 4A \rightarrow 0$ model in 2 dimensions

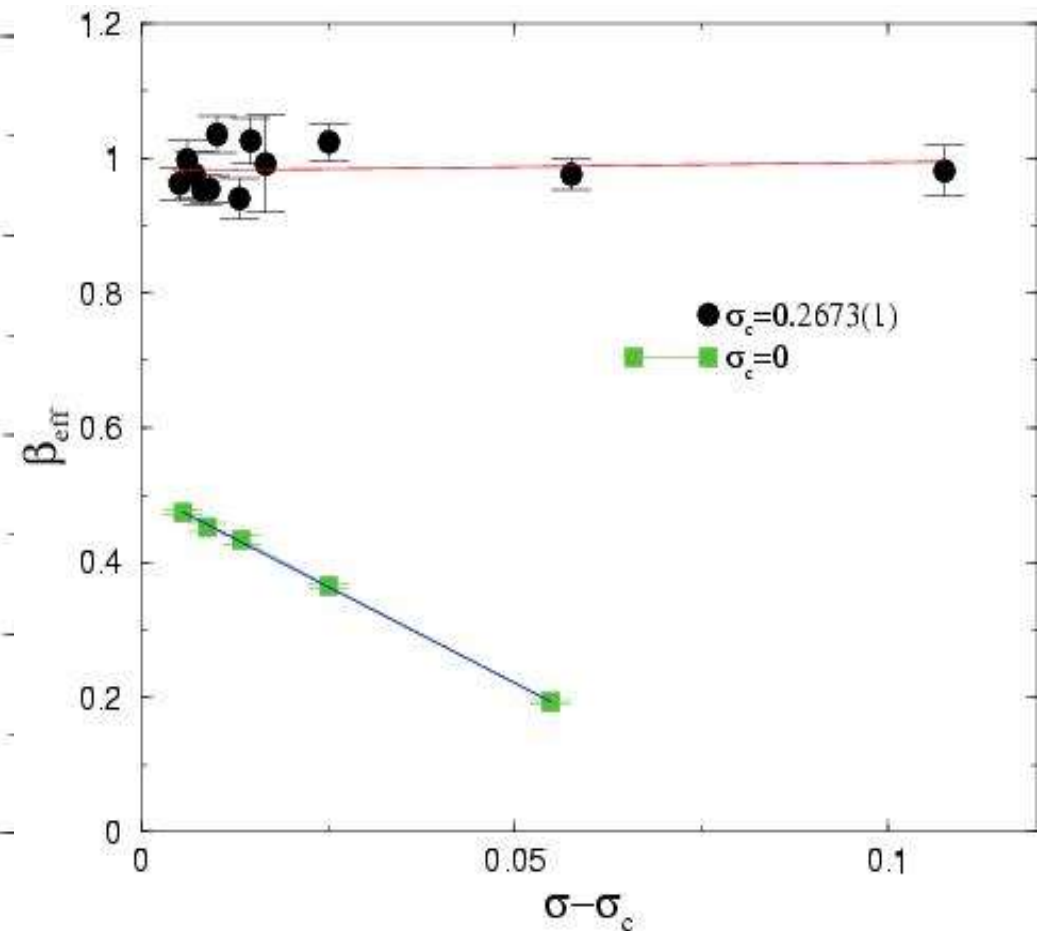
Density decay local exponents with :
 $\alpha = 0.50(1)$ (MF-PCPD) at $\sigma_c = 0.26715$

$L = 7000, t_{max} = 2 \times 10^6$ MCS, $D = 0.05$



Steady state local exponents at $D = 0.05$
 with $\beta = 0.98(2)$ (MF-PCPD).

For $D=0.9$: only mean-field with $\beta=1/2$.
 $2A \rightarrow 0$ process becomes irrelevant !



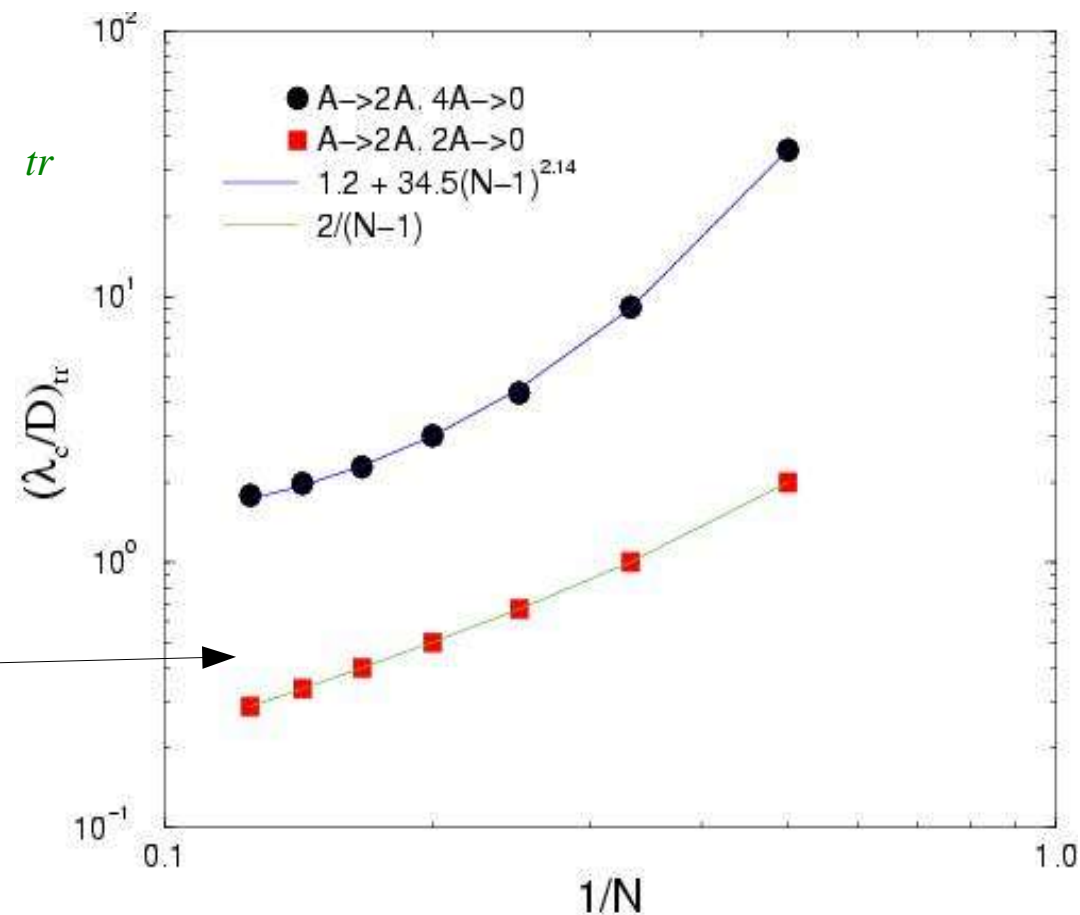
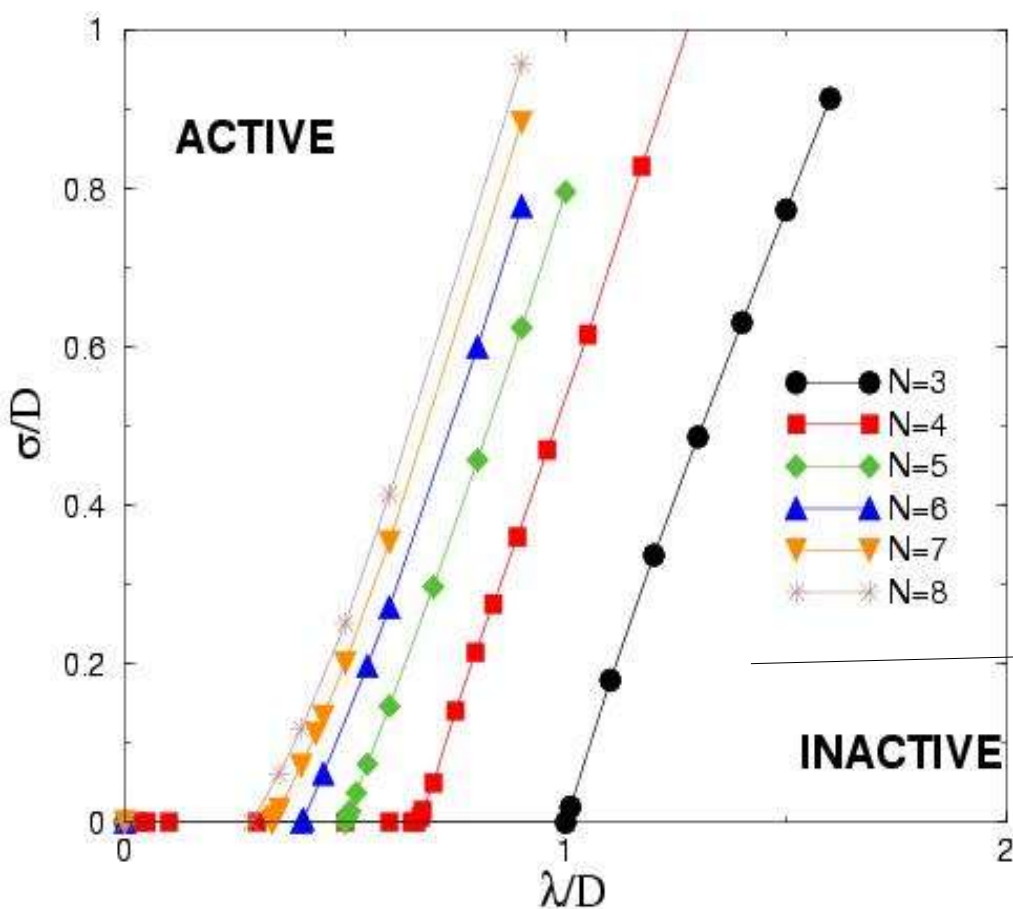
Cluster mean-field for BARW(unary)

1) $A \rightarrow 2A, 2A \rightarrow 0$: (Canet, Chaté, Delamotte, cond-mat/0403432, NPRG)

Exact relation: $(\lambda_c/D)_{tr} = 2/(N-1)$ found. For 1d: $(\lambda_c/D)_{tr} \rightarrow 0$ as $N \rightarrow \infty$.

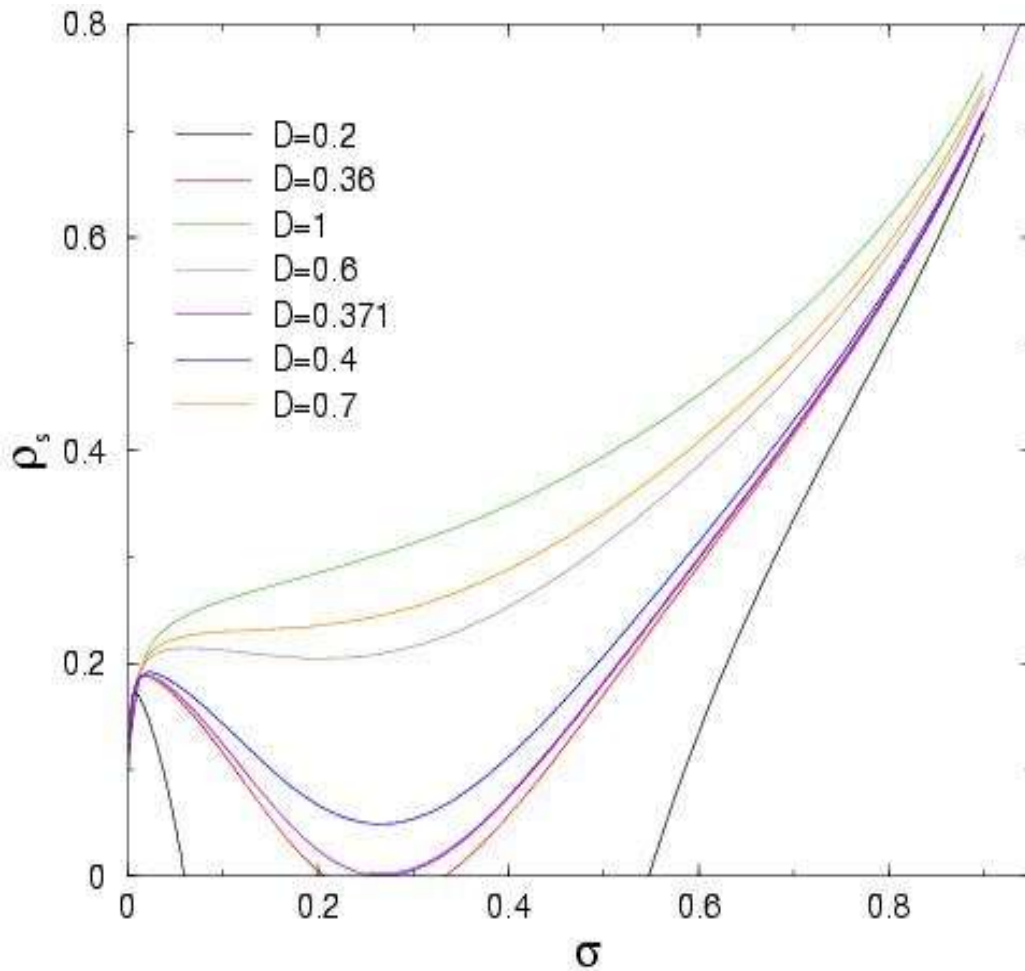
2) $A \rightarrow 2A, 4A \rightarrow 0$: $(\lambda_c/D)_{tr} \rightarrow \sim 1.2$ as $N \rightarrow \infty$.

$$MF : m > n : \beta = 1/(m-n), \alpha = 1/(m-1), \sigma_c = 0$$



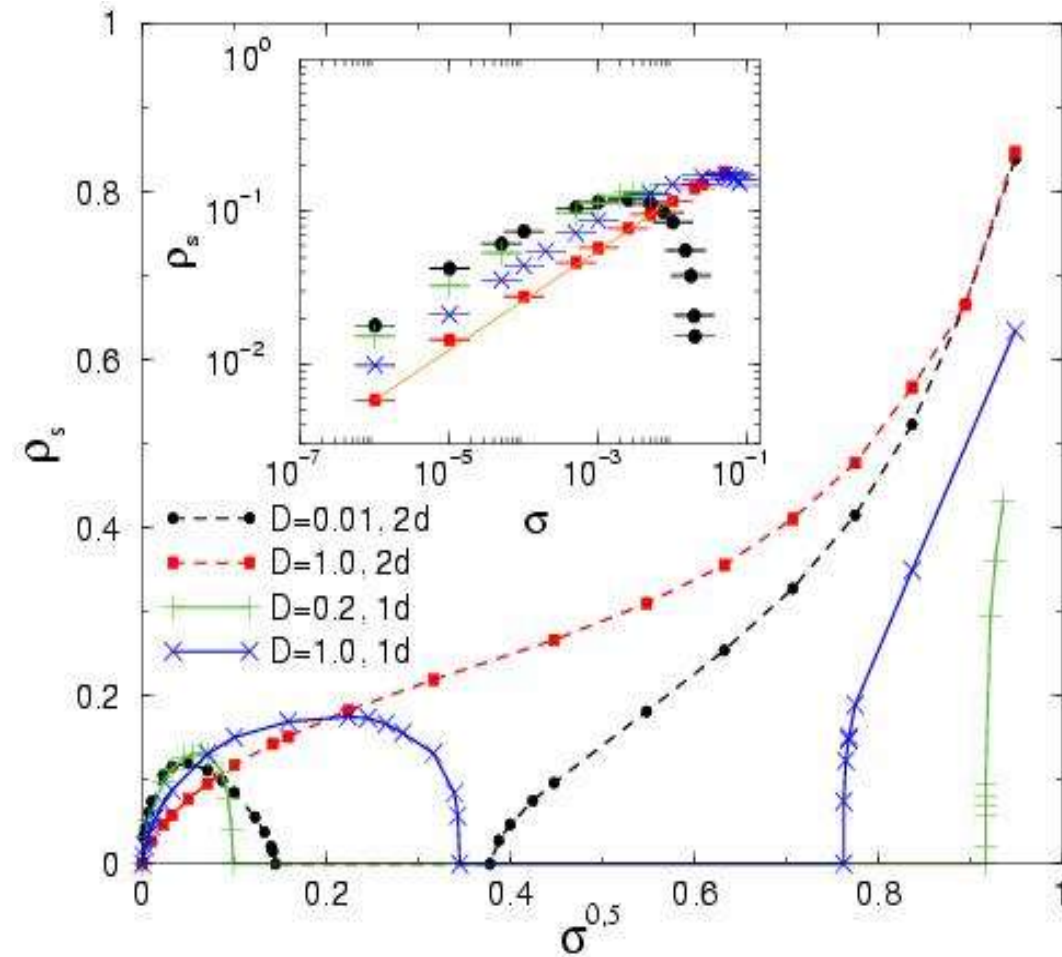
A- \rightarrow 2A, 4A- \rightarrow 0: Reentrant phase diagram

N-cluster approximation



Simulations:

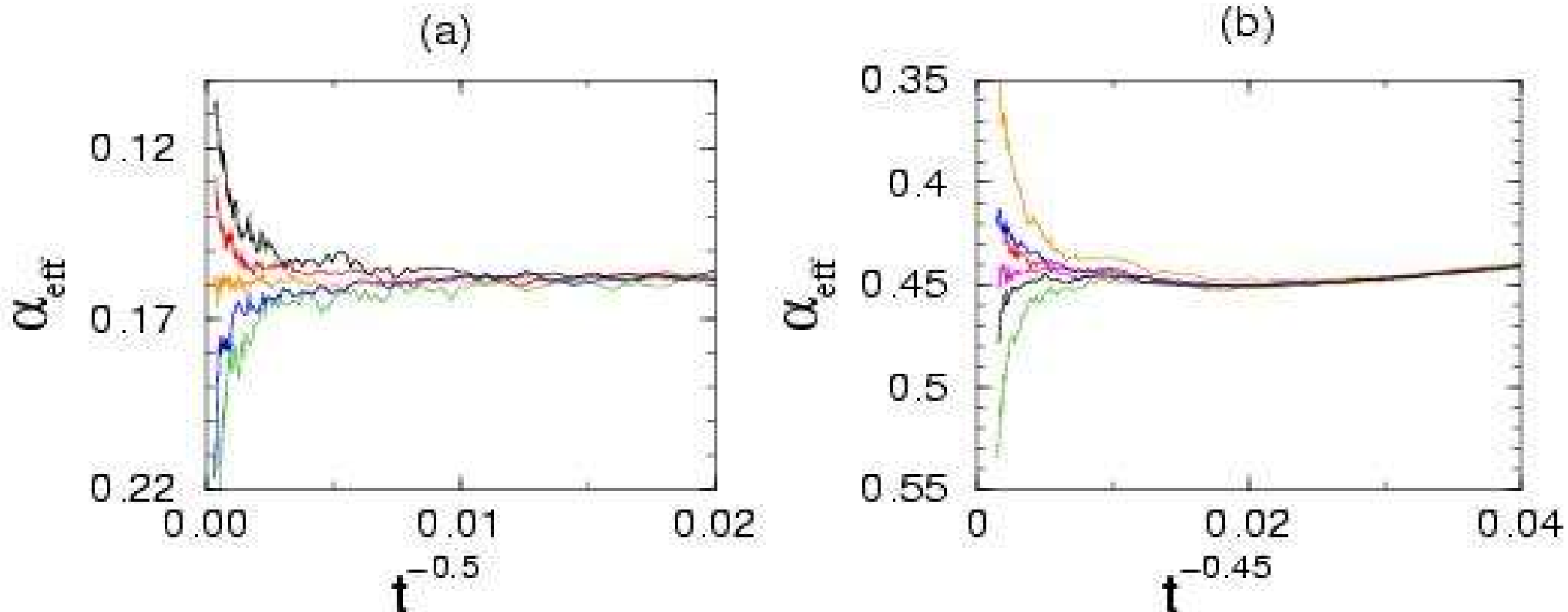
2d: $L=4000 \times 4000$, 1d: $L=100000$



Critical point decay at $\sigma_c > 0$ in (a): one and (b): two dimensions

DP class value in 1d: $\alpha = 0.159464(6)$ (I. Jensen)

2d: $\alpha = 0.4505(10)$ (Voigt and Ziff)



For low diffusions the : $A \rightarrow 2A \rightarrow 3A \rightarrow 4A \rightarrow 0$ process becomes relevant !

Summary

- In general (site restricted) reaction-diffusion systems the mean-field universality class is determined by the number of reacting particles (n,m)
- Diffusion may become relevant in case of competing reactions by changing the phase diagram and introducing nontrivial fixed points
- Diffusion dependence has been found in similar systems:
 - *R. Dickman: A \rightarrow 2A, 3A \rightarrow 0, PRA42, 6985 (1990)*
 - *M. Paessens, G. Schütz : bosonic PCPD, JPA37, 4709 (2004)*
 - *N. Menyhárd, G. Ódor : NEKIM-A, PRE68, 056106 (2003)*
- **Perturbative RG results should be revised !**
- References : *G. Ódor: PRE 67, 056114 (2003)*
G. Ódor: PRE 67, 016111 (2003)
G. Ódor: PRE 69, 036112 (2004)
G. Ódor: Rev. Mod. Phys. 76 (2004), cond-mat/0205644
G. Ódor: cond-mat/0403562, to appear in PRE
G. Ódor: cond-mat/0406247